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In a previous article (*Int. Math. Res. Not.* **2014**:10 (2014), 2730–2745) T. Orponen and the authors proved that the Fourier dimension of the graph of any real-valued function on \mathbb{R} is bounded above by 1. This partially answered a question of Kahane (1993) by showing that the graph of the Wiener process W_t (Brownian motion) is almost surely not a Salem set. In this article we complement this result by showing that the Fourier dimension of the graph of W_t is almost surely 1. In the proof we introduce a method based on Itô calculus to estimate Fourier transforms by reformulating the question in the language of Itô drift-diffusion processes and combine it with the classical work of Kahane on Brownian images.

1. Introduction and results

1A. Geometric properties of Brownian motion. Gaussian processes are standard models in modern probability theory and perhaps the most well-studied example is the *Wiener process* (or standard *Brownian motion*) $W = W_t : \mathbb{R}_{\geq 0} \rightarrow \mathbb{R}$ characterised by the properties $W_0 = 0$, the map $t \mapsto W_t$ is almost surely continuous, and W_t has independent increments such that $W_t - W_s$ for $t > s$ is normally distributed:

$$W_t - W_s \sim N(0, t - s).$$

The Wiener process has far-reaching importance throughout mathematics and it is a topic of particular interest to understand its geometric structure. This can be achieved by studying several random fractals associated to the process such as *images* $W(K) := \{W_t : t \in K\}$ of compact sets $K \subset [0, \infty)$, *level sets* $L_c(W) := \{t \in \mathbb{R} : W_t = c\}$ for $c \in \mathbb{R}$, *graphs* $G(W) := \{(t, W_t) : t \in \mathbb{R}\}$ (see [Figure 1](#)) and other more delicate constructions such as SLE_κ -curves.

The basic properties of Brownian motion mean that these random fractals enjoy a certain “statistical self-similarity”, which facilitates computation of their Hausdorff dimensions \dim_{H} . Classical results include McKean’s proof [\[1955\]](#) that $\dim_{\text{H}} W(K) = \min\{1, 2 \dim_{\text{H}} K\}$ almost surely for each compact $K \subset [0, \infty)$. Moreover, for the level sets, $\dim_{\text{H}} L_c(W) = \frac{1}{2}$ almost surely for $c = 0$ by [\[Taylor 1955\]](#) and for all $c \in \mathbb{R}$ by [\[Perkins 1981\]](#) conditioned on $L_c(W)$ being nonempty. For the Brownian graph $G(W)$, Taylor [\[1953\]](#) proved that $\dim_{\text{H}} G(W) = \frac{3}{2}$ almost surely and Beffara [\[2008\]](#) computed the Hausdorff dimensions of SLE_κ -curves. Moreover, Hausdorff dimensions for similar sets given by many other Gaussian processes, such as fractional Brownian motion, have been also considered; see, for example,

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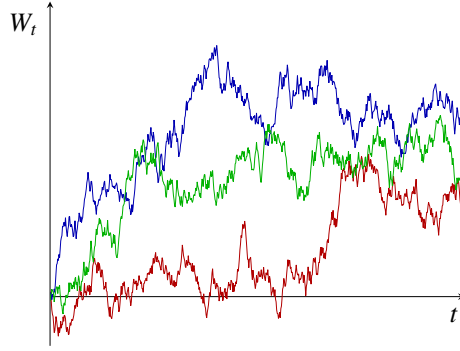


Figure 1. Three realisations of the graph $G(W)$ for the Brownian motion W_t .

Adler’s classical results [1977] for fractional Brownian graphs and the recent work [Peres and Sousi 2016] concerning variable drift.

1B. Fourier analytic properties of Brownian motion. The Hausdorff dimension is the most commonly used tool for measuring the size of a set A but there is also another fundamental notion based on Fourier analysis which reveals more arithmetic and geometric features of A , including curvature, which are not seen by the Hausdorff dimension. This is based on studying the *Fourier coefficients* of a probability measure μ on $A \subset \mathbb{R}^d$, which are defined by

$$\hat{\mu}(\xi) := \int e^{-2\pi i \xi \cdot x} d\mu(x), \quad \xi \in \mathbb{R}^d.$$

Now the size of A can be linked to the existence of probability measures μ on A with decay of Fourier coefficients $\hat{\mu}(\xi)$ when $|\xi| \rightarrow \infty$. The following connection between Hausdorff dimension and decay of Fourier coefficients is well known and goes back to Salem and Kaufman, but we refer the reader to [Mattila 2015] for the details. If $\dim_H A > s$, then A supports a probability measure μ with $|\hat{\mu}(\xi)| = O(|\xi|^{-s/2})$ “on average”, that is, $\int_{\mathbb{R}^d} |\hat{\mu}(\xi)|^2 |\xi|^{s-d} d\xi < \infty$, and vice versa the Hausdorff dimension can be bounded from below if such a measure μ can be found. It is possible, however, that $\dim_H A = s > 0$ but no measure μ on A has Fourier decay at infinity; this happens for example when A is the middle-third Cantor set in \mathbb{R} . Therefore, one defines the notion of *Fourier dimension* $\dim_F A$ of a set $A \subset \mathbb{R}^d$ as the supremum of $s \in [0, d]$ for which there exists a probability measure μ supported on A such that

$$|\hat{\mu}(\xi)| = O(|\xi|^{-s/2}) \quad \text{as } |\xi| \rightarrow \infty. \quad (1-1)$$

Then by this definition we always have $\dim_F A \leq \dim_H A$ and if the two dimensions coincide then A is called a *Salem set* or a *round set* after [Kahane 1993]. In general Fourier dimension and Hausdorff dimension have no relationship other than this; in fact, Körner [2011] established that for any $0 \leq s < t \leq 1$ it is possible to construct examples $A \subset \mathbb{R}$ with $\dim_F A = s$ and $\dim_H A = t$. Further properties of Fourier dimension were recently developed by Ekström, Persson and Schmeling [Ekström et al. 2015]. For a more in depth account of Fourier dimension, the reader is referred to [Mattila 1995; 2015].

Finding measures μ on A with polynomially decaying Fourier transform (i.e., (1-1) for some $s > 0$) has deep links to absolute continuity, arithmetic and geometric structure, and curvature. If A supports a measure μ such that (1-1) holds with $s > 1$, then Parseval's identity yields that μ is absolutely continuous to Lebesgue measure and A must contain an interval. An application of Weyl's criterion known as the Davenport–Erdős–LeVeque criterion [Davenport et al. 1963] yields that in \mathbb{R} polynomial decay of $\hat{\mu}$ guarantees that μ almost every number is normal in every base and an interesting result of Łaba and Pramanik [2009] shows that if the s in (1-1) is sufficiently close to 1 for a Frostman measure μ on $A \subset \mathbb{R}$ and there is a suitable control over the constants, see the recent work [Shmerkin 2017], then A contains nontrivial 3-term arithmetic progressions. Moreover, an analogous result also holds for higher dimensions with arithmetic patches [Chan et al. 2016].

On the curvature side, if A is a line-segment in \mathbb{R}^2 , then A cannot contain any measure with Fourier decay at infinity so A cannot be a Salem set. However, if A is an arc of a circle or more generally a 1-dimensional smooth manifold with nonvanishing curvature then the 1-dimensional Hausdorff measure μ on A satisfies (1-1) with $s = 1$; see [Mattila 2015]. In particular, A is a Salem set. In these examples of A one can observe that the important arithmetic or curvature features present are not seen from the Hausdorff dimension.

Constructing explicit Salem sets (which are not manifolds), or just sets A supporting a measure μ satisfying (1-1) for some $s > 0$, can be achieved through, for example, Diophantine approximation by [Kaufman 1980; 1981; Bluhm 1998; Queffélec and Ramaré 2003] or via thermodynamical tools by [Jordan and Sahlsten 2016]. However, for random sets it has been observed in many instances that A is either almost surely Salem or at least supports a measure μ with (1-1) for some $s > 0$. This was first done for random Cantor sets by Salem [1951], where Salem sets were also introduced. Later in his classical papers, Kahane [1966a; 1966b] found out that the Wiener process and other Gaussian processes provide natural examples.

Since Kahane and Salem, the study of Fourier analytic properties of natural sets derived from Gaussian processes and more general random fields has been an active topic. For the Brownian images, Kahane [1985b] proved that for any compact $K \subset \mathbb{R}$ the image $W(K)$ is almost surely a Salem set of Hausdorff dimension $\min\{1, 2 \dim K\}$. Kahane also established a similar result for fractional Brownian motion. Łaba and Pramanik [2009] then applied these to the additive structure of Brownian images. Later Shieh and Xiao [2006] extended Kahane's work to very general classes of Gaussian random fields. However, understanding the Fourier analytic properties of the level sets and graphs remained an important problem for some time. Kahane [1993] outlined the problem explicitly.

Problem 1.1 (Kahane). *Are the graph and level sets of a stochastic process, such as fractional Brownian motion, Salem sets?*

This precise formulation of the problem was given by Shieh and Xiao [2006, Question 2.15], but they attribute the problem to Kahane. For the Wiener process Kahane [1985a] had already established that the level sets $L_c(W)$ are Salem almost surely for any fixed $c \in \mathbb{R}$ conditioned on $L_c(W)$ being nonempty. The fractional Brownian motion case has recently been considered for $c = 0$ by Fouché and Mukeru [2013].

Kahane's problem for graphs, even in the case of the standard Brownian motion W_t , however, remained open for quite a while until, together with T. Orponen, we established that the Brownian graph $G(W)$ is

almost surely *not* a Salem set [Fraser et al. 2014]. It turned out that the reason for this is purely geometric: the proof was based on the following application of a Fourier-analytic version of Marstrand’s slicing lemma.

Theorem 1.2 [Fraser et al. 2014, Theorem 1.2]. *For any function $f : [0, 1] \rightarrow \mathbb{R}$ the Fourier dimension of the graph $G(f)$ cannot exceed 1.*

Indeed, since $\dim_{\text{H}} G(W) = \frac{3}{2} > 1$ almost surely [Taylor 1953], this answers Kahane’s problem in the negative for the Wiener process. Note that this also gives a negative answer for fractional Brownian motion since the Hausdorff dimension in that case is also strictly larger than 1 almost surely.

The methods in [Fraser et al. 2014] are purely geometric and involve no stochastic properties of Brownian motion. They also do not shed any light on the precise value for the Fourier dimension of $G(W)$. Note that even though $\dim_{\text{H}} G(f) \geq 1$ for any continuous $f : [0, 1] \rightarrow \mathbb{R}$, the Fourier dimension of a graph may take any value in the interval $[0, 1]$; see [loc. cit.]. For example, $\dim_{\text{F}} G(f) = 0$ if f is affine and, moreover, $\dim_{\text{F}} G(f) = 0$ for the *Baire generic* $f \in C[0, 1]$; see [loc. cit., Theorem 1.3].

The main result of this paper is to complete the work initiated by Kahane’s problem in the case of Brownian motion by establishing the precise almost sure value of the Fourier dimension of $G(W)$.

Theorem 1.3. *The graph $G(W)$ has Fourier dimension 1 almost surely.*

Moreover, the random measure μ we use to realise the Fourier dimension is Lebesgue measure dt on $[0, 1]$ lifted onto the graph $G(W)$ via the mapping $t \mapsto (t, W_t)$. The precise estimate we obtain is that almost surely

$$|\hat{\mu}(\xi)| = O(|\xi|^{-1/2} \sqrt{\log |\xi|}) \quad \text{as } |\xi| \rightarrow \infty, \quad (1-2)$$

which combined with Theorem 1.2 yields Theorem 1.3.

A natural direction in which to continue this line of research would be to study other Gaussian processes with different covariance structure, such as the fractional Brownian motion.

1C. Methods: Itô calculus and reduction to Brownian images. The key method we introduce to estimate the Fourier transform of the graph measure μ is based on *Itô calculus*, which has previously been a natural framework in the theory of stochastic differential equations. As far as we know, Itô calculus has not been previously considered in this Fourier analytic context. Here we discuss this method and give a brief summary of the main steps in the proof. When written in polar coordinates, (1-2) asks about the rate of decay for the integral

$$\hat{\mu}(\xi) = \int_0^1 \exp(-2\pi i u(t \cos \theta + W_t \sin \theta)) dt$$

for $\xi = u(\cos \theta, \sin \theta) \in \mathbb{R}^2$, $u > 0$, $\theta \in [0, 2\pi)$, as $u \rightarrow \infty$. There are two distinct cases we will consider depending on the direction of ξ , which we give a heuristic description of here.

If we ignore the random component $W_t \sin \theta$, that is, set $\theta = 0$ or π , then standard integration using the chain rule shows that $\hat{\mu}(\xi)$ equals the Fourier transform of Lebesgue measure dt at u , which decays to 0 with the polynomial rate $u^{-1} = |\xi|^{-1}$, so we are done for these directions. However, if θ is not equal to 0

or π , we still have a small random (nonsmooth) term $W_t \sin \theta$, so a classical change of variable formula or other tools from classical analysis cannot be used.

The key observation is that we can write $\hat{\mu}(\xi) = \int \exp(i X_t) dt$, where the stochastic process X_t satisfies the stochastic differential equation

$$dX_t := b dt + \sigma dW_t,$$

identifying it as a so-called *Itô drift-diffusion process*, where $b := -2\pi u \cos \theta$ is the *drift coefficient* of X_t and $\sigma := -2\pi u \sin \theta$ is the *diffusion coefficient* of X_t . Such processes have many useful analytic tools from Itô calculus (see [Section 2](#)) associated to them, in particular *Itô's lemma*, which works as an analogue for the chain rule. The price we pay is that Itô's lemma introduces some multiplicative error terms involving stochastic integrals, but they can be estimated with other tools from Itô calculus using moment analysis.

The estimates we obtain from Itô calculus allow us to obtain the correct Fourier decay (1-2) for μ when θ is close to 0 or π with respect to u^{-1} (more precisely, $|\sin \theta| < u^{-1/2}$), in other words, when ξ is close to pointing in the horizontal directions. Thus another estimate is needed for θ bounded away from 0 and π . This is where the classical work [\[Kahane 1985b\]](#) on Brownian images comes into play. If we completely ignore the deterministic component $t \cos \theta$, by setting $\theta = \frac{\pi}{2}$ or $\frac{3\pi}{2}$, then $\hat{\mu}(\xi)$ is the Fourier transform of the Brownian image measure ν , that is, the $t \mapsto W_t$ push-forward of the Lebesgue measure dt on $[0, 1]$ at u . Kahane [\[1985b\]](#) in fact already established that the decay of $|\hat{\nu}(u)|$ is almost surely of the order $u^{-1} \sqrt{\log u} = |\xi|^{-1} \sqrt{\log |\xi|}$ so (1-2) holds for these directions. A modification of Kahane's argument reveals that whenever $\theta \neq 0$ or π , then almost surely

$$|\hat{\mu}(\xi)| = O(|\sin \theta|^{-1} |\xi|^{-1} \sqrt{\log |\xi|});$$

see the discussion in [Section 3C](#). Now one notices that when θ approaches 0 or π , this estimate blows up, and so one cannot obtain a uniform estimate over all directions from this. However, this gives (1-2) if $|\sin \theta| \geq u^{-1/2}$, so combining with the estimates we obtained through Itô calculus, we are done. See [Section 3](#) for more details on the main steps of the proof.

1D. Other measures on the Brownian graph. [Theorem 1.3](#) and (1-2) give Fourier decay for the push-forward of the Lebesgue measure on $[0, 1]$ onto the graph $G(W)$. It would be an interesting problem to see if one can have similar results for other, possibly fractal, measures on $[0, 1]$. A possible problem could be:

Problem 1.4. *Classify measures τ on $[0, 1]$ such that for some $0 < s \leq 1$ we have*

$$|\hat{\tau}(\xi)| = O(|\xi|^{-s/2}), \quad |\xi| \rightarrow \infty,$$

and their lift μ_τ onto the graph of $G(W)$ under $t \mapsto (t, W_t)$ satisfies

$$|\hat{\mu}_\tau(\xi)| = O(|\xi|^{-s'/2}), \quad |\xi| \rightarrow \infty$$

for any $s' < s$.

This is motivated by the fact that in [\[Kahane 1985b\]](#) it is possible to transfer information on the Fourier decay (or Frostman properties) of τ onto the image measure. Thus for directions θ bounded away from 0

and π we could still bound $\hat{\mu}_\tau(\xi)$ using Kahane's work. The main problem in generalising our approach to fractal measures τ on $[0, 1]$ comes from the lack of an appropriate analogue of Itô calculus.

1E. Organisation of the paper. In [Section 2](#) we give the necessary background from Itô calculus. In [Section 3](#) we will give the proof of our main result [Theorem 1.3](#). The key estimates are obtained in [Section 3B](#) and [Section 3C](#), corresponding to the two cases discussed above.

2. Itô calculus

2A. Stochastic integration. In the proof of the main result [Theorem 1.3](#), we end up studying integrals of the form $\int f(X_t) dt$ for some stochastic processes X_t and smooth scalar functions f . As standard analysis methods cannot be applied to these integrals, we need theory from stochastic analysis. Stochastic analysis provides a pleasant framework to deal with nonsmooth processes, such as the Wiener process W_t , and still preserves many of the classical features present in the smooth setting. In this section we discuss the specific tools from *Itô calculus* which we will rely on. The main references for this section are given in the book [\[Karatzas and Shreve 1991\]](#).

Let $(\Omega, \mathcal{F}, (\mathcal{F}_t)_{t \geq 0}, \mathbb{P})$ be a filtered probability space; that is, $\mathcal{F}_t \subset \mathcal{F}$ is an increasing filtration in t . Let $W = W_t$ be the Wiener process adapted to this filtered probability space; that is, W_t is \mathcal{F}_t measurable and for each $t, s \geq 0$ the increment $W_{t+s} - W_t$ is independent of \mathcal{F}_t . We say that an \mathbb{R} - or \mathbb{C} -valued stochastic process Z_t is *adapted* if it is \mathcal{F}_t measurable for all $t \geq 0$. We will say that a real- or complex-valued adapted process Z_t is *W_t -integrable* if the *quadratic variation* $\int_0^T |Z_t|^2 dt$ is finite for any time $T \geq 0$. Given a real-valued adapted W_t -integrable stochastic process X_t , we have \mathbb{P} almost surely for any time $T \geq 0$ it is possible to construct a *stochastic integral*

$$\int_0^T X_t dW_t$$

of X_t with respect to W_t in the sense of Itô; see [\[Karatzas and Shreve 1991, Chapter 3.2\]](#). We use the differential notation $dU_t = X_t dW_t$ to mean that \mathbb{P} almost surely $U_T - U_0$ is the stochastic integral of X_t with respect to W_t at time $T \geq 0$.

We mainly deal with complex-valued stochastic processes, so for the sake of convenience we will also introduce the *complex-valued* stochastic integral for a \mathbb{C} -valued W_t -integrable adapted process Z_t , defined coordinatewise using real integrals:

$$\int_0^T Z_t dW_t := \int_0^T \operatorname{Re} Z_t dW_t + i \int_0^T \operatorname{Im} Z_t dW_t,$$

where the real integrals are standard \mathbb{R} -valued stochastic integrals with respect to the Wiener process W_t . We write $dZ_t := dX_t + i dY_t$ for a complex-valued process $Z_t = X_t + iY_t$ with \mathbb{R} -valued X_t and Y_t .

2B. Itô drift-diffusion processes. The main class of adapted processes to which we apply Itô calculus is given by Wiener processes with drift and diffusion coefficients. These are called Itô drift-diffusions:

Definition 2.1 (Itô drift-diffusion process). A real- or complex-valued adapted stochastic process X_t is called an *Itô drift-diffusion process* if there exists a Lebesgue integrable adapted b_t and W_t -integrable

adapted σ_t such that X_t satisfies the stochastic differential equation

$$dX_t = b_t dt + \sigma_t dW_t.$$

For Itô drift-diffusion processes there exists the following important analogue of the change of variable formula, which follows from robustness of Taylor expansions for stochastic differentials:

Lemma 2.2 (Itô's lemma). *Let X_t be an Itô drift-diffusion process and $f : \mathbb{R} \rightarrow \mathbb{R}$ twice differentiable. Then $f(X_t)$ is an Itô drift-diffusion process such that \mathbb{P} almost surely for any $T \geq 0$ we have*

$$f(X_T) - f(X_0) = \int_0^T (b_t f'(X_t) + \frac{1}{2} \sigma_t^2 f''(X_t)) dt + \int_0^T \sigma_t f'(X_t) dW_t.$$

Itô's lemma was given in this pathwise form in [Karatzas and Shreve 1991, Theorem 3.3]. By using the definition of the complex-valued stochastic integral, we can also obtain a complex-valued Itô's lemma:

Lemma 2.3 (complex Itô's lemma). *Let X_t be an Itô drift-diffusion process and $f : \mathbb{R} \rightarrow \mathbb{C}$ twice differentiable. Then $f(X_t)$ is an Itô drift-diffusion process such that for \mathbb{P} almost surely for any $T \geq 0$ we have*

$$f(X_T) - f(X_0) = \int_0^T (b_t f'(X_t) + \frac{1}{2} \sigma_t^2 f''(X_t)) dt + \int_0^T \sigma_t f'(X_t) dW_t.$$

Proof. We can write $f = f_1 + if_2$ for real-valued twice differentiable $f_1, f_2 : \mathbb{R} \rightarrow \mathbb{R}$. Then the derivatives satisfy $f' = f'_1 + if'_2$ and $f'' = f''_1 + if''_2$. Moreover, by Itô's lemma (Lemma 2.2) we obtain for each $j = 1, 2$ that

$$df_j(X_t) = (b_t f'_j(X_t) + \frac{1}{2} \sigma_t^2 f''_j(X_t)) dt + \sigma_t f'_j(X_t) dW_t.$$

Then by the convention $df(X_t) = df_1(X_t) + idf_2(X_t)$ this gives

$$df(X_t) = (b_t f'(X_t) + \frac{1}{2} \sigma_t^2 f''(X_t)) dt + \sigma_t f'(X_t) dW_t$$

as required. □

2C. Moment estimation. Itô's lemma allows us to pass from integrals of the form $\int_0^T f(X_t) dt$ to $\int_0^T g(X_t) dW_t$ for functions g obtained from derivatives of f . In our case we will end up trying to understand the higher moments of the stochastic integrals $\int_0^T g(X_t) dW_t$, which will tell us about the distribution of these integrals. A very standard tool to compute the moments in Itô calculus are the *Itô isometry* and more general *Burkholder–Davis–Gundy inequalities* [Burkholder et al. 1972], which allow us to pass from stochastic integrals to their quadratic variations (that just involve Lebesgue integral).

Lemma 2.4 (Burkholder–Davis–Gundy inequality). *Let X_t be a real-valued W_t -integrable adapted process. Then for all $1 \leq p < \infty$ we have*

$$\mathbb{E} \left[\left(\sup_{0 \leq s \leq 1} \left| \int_0^s X_t dW_t \right| \right)^{2p} \right] \leq 2\sqrt{10p} \mathbb{E} \left[\left(\int_0^1 X_t^2 dt \right)^p \right].$$

This version with the constant $2\sqrt{10p}$ was given by Peškir [1996].

3. Proof of the main result

3A. Preliminaries and overview of the proof. Let us now review how we will prove (1-2) and thus Theorem 1.3. Fix $\xi = u(\cos \theta, \sin \theta) \in \mathbb{R}^2$ with modulus $u > 0$ and argument $\theta \in [0, 2\pi)$. Notice that by the definition of the graph measure μ , the Fourier transform has the form

$$\hat{\mu}(\xi) = \int_0^1 \exp(i X_t) dt,$$

where X_t is the real-valued stochastic process

$$X_t := -2\pi u(t \cos \theta + W_t \sin \theta). \quad (3-1)$$

The first observation is that X_t is an adapted W_t -integrable process and in fact an Itô drift-diffusion process (recall Definition 2.1) satisfying

$$dX_t = b dt + \sigma dW_t$$

for deterministic and time independent coefficients $b = -2\pi u \cos \theta$ and $\sigma = -2\pi u \sin \theta$. The proof of bounding $\hat{\mu}(\xi)$ will heavily depend on the value of the angle θ we have for ξ and in particular how close the determining angle θ is to 0, π or 2π with respect to $u^{-1/2}$. For this purpose, we define the notions of *horizontal* and *vertical* angles:

Definition 3.1 (horizontal and vertical angles). Define the threshold angle

$$\theta_u := \min\{u^{-1/2}, \frac{\pi}{4}\}.$$

Partition the angles $[0, 2\pi)$ using θ_u into the *horizontal angles*

$$H_u := [0, \theta_u] \cup [\pi - \theta_u, \pi + \theta_u] \cup [2\pi - \theta_u, 2\pi)$$

and the *vertical angles*

$$V_u := [0, 2\pi) \setminus H_u.$$

In other words H_u contains the θ_u neighborhoods of 0 and π on the circle mod 2π and V_u the $\frac{\pi}{2} - \theta_u$ neighborhoods of $\frac{\pi}{2}$ and $\frac{3\pi}{2}$ respectively; see Figure 2.

The proof will split into two cases in Sections 3B and 3C for bounding the Fourier transform $\hat{\mu}(\xi)$ depending on whether $\theta \in H_u$ or $\theta \in V_u$:

- (1) Section 3B concerns angles $\theta \in H_u$, that is, close to horizontal directions 0 or π , and as mentioned in the Introduction our main hope here is that the smallness (with respect to $u^{-1/2}$) of the diffusion component bW_t will help us in transferring the decay of Lebesgue measure to the decay of $\hat{\mu}$. This is where Itô's lemma (see Lemma 2.3) becomes crucial as it can be applied to the process $f(X_t)$ with the function $f(x) = \exp(ix)$.
- (2) Section 3C handles the angles $\theta \in V_u$ and here the plan is to use the fact that we are $u^{-1/2}$ -bounded away from horizontal angles to ignore the drift component bt of the drift-diffusion process X_t and apply Kahane's bound for these directions. This turns out to be possible due to a representation of the higher moments Kahane obtained in his result on Brownian images.

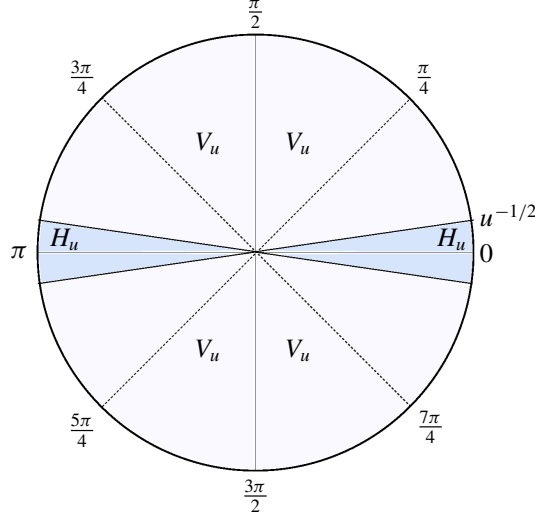


Figure 2. Splitting of $[0, 2\pi)$ to horizontal angles H_u and the vertical angles V_u .

It turns out that in both Sections 3B and 3C we only obtain decay of the Fourier transform $\hat{\mu}(k)$ for k in an ε -grid $\varepsilon\mathbb{Z}^2$ for all small $\varepsilon > 0$. Here the randomness will depend on $\varepsilon > 0$ but thanks to an argument also used by Kahane [1985b], one can pass from this information to the full decay almost surely. See Section 3D for the details.

Let us now proceed to bound $|\hat{\mu}(\xi)|$. In both Sections 3B and 3C below we will end up bounding trigonometric functions with respect to θ_u and for this purpose we will need the following standard bounds, which we record here for convenience:

Lemma 3.2 (trigonometric bounds). *We have the following bounds:*

(1) If $\theta \in H_u$, then

$$|\sin \theta| \leq u^{-1/2} \quad \text{and} \quad |\cos \theta| \geq \frac{1}{\sqrt{2}}.$$

(2) If $\theta \in V_u$, then

$$|\sin \theta| \geq \min \left\{ \frac{2}{\pi} u^{-1/2}, \frac{1}{\sqrt{2}} \right\}.$$

Proof. For $\alpha \in [0, \frac{\pi}{2}]$ we have that both cosine and sine are nonnegative. Moreover, here $\frac{2}{\pi}\alpha \leq \sin \alpha \leq \alpha$. Thus for $\theta \in [0, \theta_u]$ we have

$$\sin \theta \leq \theta \leq \theta_u \leq u^{-1/2} \quad \text{and} \quad \cos \theta \geq \cos \theta_u \geq \cos \frac{\pi}{4} = \frac{1}{\sqrt{2}}.$$

and for $\theta \in (\theta_u, \frac{\pi}{2}]$ as $\sin \frac{\pi}{4} = \frac{1}{\sqrt{2}}$ we obtain

$$\sin \theta \geq \min \left\{ \frac{2}{\pi} u^{-1/2}, \frac{1}{\sqrt{2}} \right\}.$$

This gives the claim as we may reduce the estimates back to the estimates for $\theta \in [0, \frac{\pi}{2}]$ by using standard invariance identities for sine and cosine. \square

3B. Horizontal angles. When $\theta \in H_u$ we will first obtain the following estimate on ε -grids:

Lemma 3.3. *Fix $\varepsilon > 0$. Almost surely there exists a random constant $C_\omega > 0$ such that for any $\mathbf{k} = u(\cos \theta, \sin \theta) \in \varepsilon\mathbb{Z}^2 \setminus \{0\}$ with $\theta \in H_u$ we have*

$$|\hat{\mu}(\mathbf{k})| \leq C_\omega |\mathbf{k}|^{-1/2}.$$

Given $\xi = u(\cos \theta, \sin \theta) \in \mathbb{R}^2 \setminus \{0\}$ and a realisation (W_t) , define a random variable $T = T_\omega(\xi) \in [0, 1]$, to be the minimum value of $t \in [0, 1]$ such that

$$X_t = \begin{cases} -2\pi \lceil u(\cos \theta + W_1 \sin \theta) \rceil & \text{if } X_1 \geq 0, \\ -2\pi \lfloor u(\cos \theta + W_1 \sin \theta) \rfloor & \text{if } X_1 < 0. \end{cases}$$

Such a time T exists almost surely since $X_0 = 0$ and X_t is almost surely continuous (since W_t is almost surely continuous). Splitting the integral of Z_t up into “complete rotations” and “what is left over”, one obtains

$$\int_0^1 Z_t dt = \int_0^T Z_t dt + \int_T^1 Z_t dt.$$

For the integral over $[T, 1]$ we get the following estimate.

Lemma 3.4. *Almost surely there exists a random constant $C_\omega > 0$ such that for any $\xi = u(\cos \theta, \sin \theta) \in \mathbb{R}^2 \setminus \{0\}$ with $\theta \in H_u$ we have*

$$\left| \int_T^1 Z_t dt \right| \leq C_\omega |\xi|^{-1/2}.$$

Proof. Since W_t is almost surely continuous, there almost surely exists a random constant $M_\omega > 1$ such that $W_t \in [-M_\omega, M_\omega]$ for all $t \in [0, 1]$. Define the real-valued process

$$Y_t := u(t \cos \theta + W_t \sin \theta),$$

so $X_t = -2\pi Y_t$. Suppose $X_1 \geq 0$. In this case $Y_T = \lceil Y_1 \rceil \leq 0$ and so $Y_1 + 1 \geq Y_T \geq Y_1$. Moreover, when $X_1 < 0$ we have $Y_T = \lfloor Y_1 \rfloor > 0$ and $Y_1 \geq Y_T \geq Y_1 - 1$. Thus no matter what the sign of X_1 is, we always have almost surely

$$u(\cos \theta + W_1 \sin \theta) + 1 \geq u(T \cos \theta + W_T \sin \theta) \geq u(\cos \theta + W_1 \sin \theta) - 1.$$

Therefore, in the case $\cos \theta > 0$ we obtain

$$T \geq 1 + W_1 \frac{\sin \theta}{\cos \theta} - W_T \frac{\sin \theta}{\cos \theta} - \frac{1}{u \cos \theta},$$

and when $\cos \theta < 0$ we have

$$T \geq 1 + W_1 \frac{\sin \theta}{\cos \theta} - W_T \frac{\sin \theta}{\cos \theta} + \frac{1}{u \cos \theta}.$$

Since $u \in H_u$, [Lemma 3.2](#) together with $W_t \in [-M_\omega, M_\omega]$ yields

$$T \geq 1 - 2\sqrt{2}M_\omega u^{-1/2} - \frac{\sqrt{2}}{u}.$$

Recalling $M_\omega > 1$ this gives

$$\left| \int_T^1 Z_t dt \right| \leq \int_T^1 |Z_t| dt = 1 - T \leq 4M_\omega u^{-1/2}$$

as required. \square

We now estimate the integral over $[0, T]$, which is where Itô calculus comes into play.

Lemma 3.5. *Fix $\varepsilon > 0$. Almost surely there exists a random constant $C_\omega > 0$ such that for any $\mathbf{k} = u(\cos \theta, \sin \theta) \in \varepsilon \mathbb{Z}^2 \setminus \{0\}$ with $\theta \in H_u$ we have*

$$\left| \int_0^T Z_t dt \right| \leq C_\omega |\mathbf{k}|^{-1/2}.$$

To prove [Lemma 3.5](#), we first need to compute the higher-order moments of the random variable $\int_0^T Z_t dt$.

Lemma 3.6. *For any $p \in \mathbb{N}$ and $\xi = u(\cos \theta, \sin \theta) \in \mathbb{R}^2 \setminus \{0\}$ with $\theta \in H_u$, the $(2p)$ -th moment satisfies*

$$\mathbb{E} \left| \int_0^T Z_t dt \right|^{2p} \leq 13p^{1/2} 4^p |\xi|^{-p}.$$

Proof. Recall that

$$X_t = -2\pi u(t \cos \theta + W_t \sin \theta)$$

is an Itô drift-diffusion process satisfying the stochastic differential equation

$$dX_t = b dt + \sigma dW_t$$

for deterministic and time-independent coefficients $b = -2\pi u \cos \theta$ and $\sigma = -2\pi u \sin \theta$. Writing

$$f(x) := \exp(ix), \quad x \in \mathbb{R},$$

we have $Z_t = f(X_t)$, $f'(x) = i \exp(ix)$ and $f''(x) = -\exp(ix)$. Thus by complex Itô's lemma (see [Lemma 2.3](#)) we have \mathbb{P} almost surely

$$f(X_T) - f(X_0) = (bi - \sigma^2/2) \int_0^T f(X_t) dt + \sigma i \int_0^T f(X_t) dW_t. \quad (3-2)$$

Note that $T_\omega \leq 1$ is random and only \mathcal{F}_1 measurable; thus it is not a stopping time. However, as [Lemma 2.3](#) is given *pathwise*, that is, \mathbb{P} almost surely Itô's lemma holds for any time $T \geq 0$, then as T_ω is \mathbb{P} almost surely well-defined, we have (3-2) almost surely. Since X_0 and X_T are 2π multiples of integers by definition, we have $f(X_T) = f(X_0) = 1$. Thus (3-2) gives

$$\int_0^T f(X_t) dt = -\frac{\sigma i}{bi - \sigma^2/2} \int_0^T f(X_t) dW_t$$

Since b and σ are deterministic, this yields that the $(2p)$ -th moment satisfies

$$\mathbb{E} \left| \int_0^T f(X_t) dt \right|^{2p} = \left| \frac{\sigma i}{bi - \sigma^2/2} \right|^{2p} \mathbb{E} \left| \int_0^T f(X_t) dW_t \right|^{2p}.$$

Applying the Burkholder–Davis–Gundy inequality (see [Lemma 2.4](#)) for the process $\cos X_t$ gives

$$\mathbb{E} \left[\left| \int_0^T \cos X_t dW_t \right|^{2p} \right] \leq \mathbb{E} \left[\left(\sup_{0 \leq s \leq 1} \left| \int_0^s \cos X_t dW_t \right| \right)^{2p} \right] \leq 2\sqrt{10p} \mathbb{E} \left(\int_0^1 \cos^2 X_t dt \right)^p \leq 2\sqrt{10p}$$

since $\cos^2 \leq 1$. Similar application for the process $\sin X_t$ gives

$$\mathbb{E} \left[\left| \int_0^T \sin X_t dW_t \right|^{2p} \right] \leq 2\sqrt{10p}.$$

By Euler’s formula, we can write $f(X_t) = \cos X_t + i \sin X_t$ and so

$$\int_0^T f(X_t) dW_t = \int_0^T \cos X_t dW_t + i \int_0^T \sin X_t dW_t.$$

Hence

$$\begin{aligned} \mathbb{E} \left| \int_0^T f(X_t) dW_t \right|^{2p} &= \mathbb{E} \left[\left(\left| \int_0^T \cos X_t dW_t \right|^2 + \left| \int_0^T \sin X_t dW_t \right|^2 \right)^p \right] \\ &\leq \mathbb{E} \left[2^p \left| \int_0^T \cos X_t dW_t \right|^{2p} + 2^p \left| \int_0^T \sin X_t dW_t \right|^{2p} \right] \\ &= 2^p \left(\mathbb{E} \left[\left| \int_0^T \cos X_t dW_t \right|^{2p} \right] + \mathbb{E} \left[\left| \int_0^T \sin X_t dW_t \right|^{2p} \right] \right) \leq 2^p 4\sqrt{10p}. \end{aligned}$$

Moreover, as $\theta \in H_u$ we have by [Lemma 3.2](#) that $\cos^2 \theta \geq \frac{1}{2}$ and $\sin^2 \theta \leq u^{-1}$. Hence

$$\left| \frac{\sigma i}{bi - \sigma^2/2} \right|^2 = \frac{\sigma^2}{b^2 + \sigma^4/4} \leq \frac{\sigma^2}{b^2} = \frac{4\pi^2 u^2 \sin^2 \theta}{4\pi^2 u^2 \cos^2 \theta} = \frac{\sin^2 \theta}{\cos^2 \theta} \leq 2u^{-1}.$$

Therefore,

$$\mathbb{E} \left| \int_0^T f(X_t) dt \right|^{2p} \leq 4\sqrt{10p} 4^p \theta^{2p} \leq 13p^{1/2} 4^p u^{-p}.$$

as required. □

Proof of [Lemma 3.5](#). Fix $\varepsilon > 0$. Then for all $\mathbf{k} \in \varepsilon\mathbb{Z}^2 \setminus \{0\}$ define the random variable

$$I(\mathbf{k}) := \left(\int_0^T Z_t dt \right) \cdot \chi_A(\mathbf{k}),$$

where χ_A is the indicator function on the set

$$A := \{ \xi = u(\cos \theta, \sin \theta) \in \mathbb{R}^2 \setminus \{0\} : \theta \in H_u \}.$$

Note that $I(\mathbf{k})$ is well-defined and finite since $|\int_0^T Z_t dt| \leq 1$ by $|\exp(ix)| = 1$. [Lemma 3.6](#) now yields for any $\mathbf{k} \in \varepsilon\mathbb{Z}^2 \setminus \{0\}$ and $p \in \mathbb{N}$ that

$$\mathbb{E} |I(\mathbf{k})|^{2p} \leq 13p^{1/2} 4^p |\mathbf{k}|^{-p}$$

as when $\mathbf{k} \notin A$ we have $I(\mathbf{k}) \equiv 0$. Write $p_{\mathbf{k}} = \lfloor \log |\mathbf{k}| \rfloor$. Then

$$\mathbb{E} \sum_{\mathbf{k} \in \varepsilon \mathbb{Z}^2 \setminus \{0\}} |\mathbf{k}|^{-3} \frac{|I(\mathbf{k})|^{2p_{\mathbf{k}}}}{13 p_{\mathbf{k}}^{1/2} 4^{p_{\mathbf{k}}} |\mathbf{k}|^{-p_{\mathbf{k}}}} \leq \sum_{\mathbf{k} \in \varepsilon \mathbb{Z}^2 \setminus \{0\}} |\mathbf{k}|^{-3} < \infty.$$

This means that the summands tend to 0 almost surely as $|\mathbf{k}| \rightarrow \infty$ and so we can find a random constant $C_{\omega} > 0$ such that for all $\mathbf{k} \in \varepsilon \mathbb{Z}^2 \setminus \{0\}$ we have

$$|\mathbf{k}|^{-3} \frac{|I(\mathbf{k})|^{2p_{\mathbf{k}}}}{13 p_{\mathbf{k}}^{1/2} 4^{p_{\mathbf{k}}} |\mathbf{k}|^{-p_{\mathbf{k}}}} \leq C_{\omega}.$$

Therefore, by possibly making C_{ω} bigger we obtain

$$|I(\mathbf{k})| \leq C_{\omega} |\mathbf{k}|^{-1/2}.$$

This holds for each $\mathbf{k} \in \varepsilon \mathbb{Z}^2 \setminus \{0\}$, so by the definition of $I(\mathbf{k})$ we have whenever $\mathbf{k} = u(\cos \theta, \sin \theta) \in \varepsilon \mathbb{Z}^2 \setminus \{0\}$ with $\theta \in H_u$ that

$$\left| \int_0^T Z_t dt \right| \leq C_{\omega} u^{-1/2}$$

as claimed. □

We are now in position to complete the proof of [Lemma 3.3](#).

Proof of Lemma 3.3. Fix $\varepsilon > 0$. By the splitting

$$\hat{\mu}(\xi) = \int_0^1 Z_t dt = \int_0^T Z_t dt + \int_T^1 Z_t dt$$

and [Lemmas 3.4](#) and [3.5](#), we have that almost surely there exists a constant $C_{\omega} > 0$ such that for all $\mathbf{k} = u(\cos \theta, \sin \theta) \in \varepsilon \mathbb{Z}^2 \setminus \{0\}$ with $\theta \in H_u$ we have

$$|\hat{\mu}(\mathbf{k})| \leq \left| \int_0^T Z_t dt \right| + \left| \int_T^1 Z_t dt \right| \leq C_{\omega} |\mathbf{k}|^{-1/2}$$

as required. □

3C. Vertical angles. In this section we apply Kahane's work to obtain Fourier decay estimates when $\theta \in V_u$.

Lemma 3.7. Fix $\varepsilon > 0$. Almost surely there exists a random constant $C_{\omega} > 0$ such that for any $\mathbf{k} = u(\cos \theta, \sin \theta) \in \varepsilon \mathbb{Z}^2 \setminus \{0\}$ with $\theta \in V_u$ we have

$$|\hat{\mu}(\mathbf{k})| \leq C_{\omega} |\mathbf{k}|^{-1/2} \sqrt{\log |\mathbf{k}|}.$$

Let us discuss a few estimates obtained in [\[Kahane 1985b\]](#). Let ν be the push-forward of Lebesgue measure on $[0, 1]$ under the map $t \mapsto W_t$; that is, ν is the Brownian image of Lebesgue measure. Kahane established the following:

Theorem 3.8 [\[Kahane 1985b, page 255\]](#). Almost surely

$$|\hat{\nu}(v)| \leq O(|v|^{-1} \sqrt{\log |v|}) \quad \text{as } |v| \rightarrow \infty.$$

The key ingredient for the proof of [Theorem 3.8](#) was based on establishing the following bound for the higher moments:

Lemma 3.9 [[Kahane 1985b](#), page 254, estimate (2)]. *There exists a constant $C > 0$ such that for any $v \in \mathbb{R} \setminus \{0\}$ and any $p \in \mathbb{N}$ we have*

$$\mathbb{E}|\hat{v}(v)|^{2p} \leq C^p p^p |v|^{-2p}.$$

We can use [Lemma 3.9](#) to give a bound on the higher moments in our setting, but with the price that the exponent will increase from $-2p$ to $-p$.

Lemma 3.10. *There exists a constant $C > 0$ such that for any $p \in \mathbb{N}$ and $\xi = u(\cos \theta, \sin \theta) \in \mathbb{R}^2 \setminus \{0\}$ with $\theta \in V_u$ the $(2p)$ -th moment satisfies*

$$\mathbb{E}|\hat{\mu}(\xi)|^{2p} \leq C^p p^p |\xi|^{-p}.$$

Proof. Write $\mathbf{t} = (t_1, \dots, t_p) \in [0, 1]^p$ and $d\mathbf{t}$ as the Lebesgue measure on $[0, 1]^p$. Given $\mathbf{t}, \mathbf{s} \in [0, 1]^p$, we define

$$\varphi(\mathbf{t}, \mathbf{s}) := \sum_{k=1}^p (t_k - s_k), \quad \psi(\mathbf{t}, \mathbf{s}) := \sum_{k=1}^p (W_{t_k} - W_{s_k}), \quad \text{and} \quad \Psi(\mathbf{t}, \mathbf{s}) := \mathbb{E}|\varphi(\mathbf{t}, \mathbf{s})|^2.$$

By the definition of μ_θ, μ and the Fourier transform, and using the fact that the multivariate process

$$X(\mathbf{t}, \mathbf{s}) := -2\pi \cos(\theta)\varphi(\mathbf{t}, \mathbf{s}) - 2\pi \sin(\theta)\psi(\mathbf{t}, \mathbf{s})$$

is Gaussian with mean $-2\pi \cos(\theta)\varphi(\mathbf{t}, \mathbf{s})$ and variance $4\pi^2 \sin^2(\theta)\Psi(\mathbf{t}, \mathbf{s})$, we have through Fubini's theorem and the formula for the characteristic function that

$$\begin{aligned} \mathbb{E}|\hat{\mu}(\xi)|^{2p} &= \mathbb{E} \int_{[0,1]^p} \int_{[0,1]^p} \exp(-2\pi i u (\cos(\theta)\varphi(\mathbf{t}, \mathbf{s}) + \sin(\theta)\psi(\mathbf{t}, \mathbf{s}))) d\mathbf{t} d\mathbf{s} \\ &= \int_{[0,1]^p} \int_{[0,1]^p} \mathbb{E} \exp(iu X(\mathbf{t}, \mathbf{s})) d\mathbf{t} d\mathbf{s} \\ &= \int_{[0,1]^p} \int_{[0,1]^p} \exp(-2\pi i \cos(\theta)u\varphi(\mathbf{t}, \mathbf{s}) - 2\pi^2 |u \sin(\theta)|^2 \Psi(\mathbf{t}, \mathbf{s})) d\mathbf{t} d\mathbf{s}. \end{aligned}$$

Thus by taking absolute values inside the integrals, and observing that $|\exp(ix)| = 1$ for any $x \in \mathbb{R}$, we obtain

$$\mathbb{E}|\hat{\mu}(\xi)|^{2p} \leq \int_{[0,1]^p} \int_{[0,1]^p} \exp(-2\pi^2 |u \sin(\theta)|^2 \Psi(\mathbf{t}, \mathbf{s})) d\mathbf{t} d\mathbf{s}. \quad (3-3)$$

On the other hand, by doing the expansion again for the Fourier transform \hat{v} of the image measure v at $v := u \sin(\theta) \in \mathbb{R} \setminus \{0\}$ we see that

$$\mathbb{E}|\hat{v}(v)|^{2p} = \mathbb{E} \int_{[0,1]^p} \int_{[0,1]^p} \exp(-2\pi i v \psi(\mathbf{t}, \mathbf{s})) d\mathbf{t} d\mathbf{s} = \int_{[0,1]^p} \int_{[0,1]^p} \exp(-2\pi^2 v^2 \Psi(\mathbf{t}, \mathbf{s})) d\mathbf{t} d\mathbf{s},$$

which is equal to (3-3). Thus by [Lemma 3.9](#) we have

$$\mathbb{E}|\hat{\mu}(\xi)|^{2p} \leq C^p p^p |v|^{-2p}.$$

Since $\theta \in V_u$ we have $|\sin \theta| \geq \min\left\{\frac{2}{\pi}u^{-1/2}, \frac{1}{\sqrt{2}}\right\}$. When $|\sin \theta| \geq \frac{1}{\sqrt{2}}$ we obtain

$$C^p p^p |v|^{-2p} \leq (2C)^p p^p u^{-2p} \leq (2C)^p p^p u^{-p}.$$

On the other hand, if $|\sin \theta| \geq \frac{2}{\pi}u^{-1/2}$ we have

$$C^p p^p |v|^{-2p} \leq C^p p^p (2u^{-1/2}/\pi)^{-2p} u^{-2p} \leq (C\pi^2/4)^p p^p u^{-p}. \quad \square$$

Now we can complete the proof of [Lemma 3.7](#) for vertical directions:

Proof of Lemma 3.7. Fix $\varepsilon > 0$. Then for all $\mathbf{k} = u(\cos \theta, \sin \theta) \in \varepsilon\mathbb{Z}^2 \setminus \{0\}$ define the random variable

$$F(\mathbf{k}) := \hat{\mu}(\mathbf{k})\chi_B(\mathbf{k}),$$

where

$$B := \{\xi = u(\cos \theta, \sin \theta) \in \mathbb{R}^2 \setminus \{0\} : \theta \in V_u\}.$$

Now $F(\mathbf{k})$ is a well-defined finite random variable as $|\hat{\mu}(\mathbf{k})| \leq 1$ for any \mathbf{k} . From [Lemma 3.10](#) we obtain for any $\mathbf{k} \in \varepsilon\mathbb{Z}^2 \setminus \{0\}$ and $p \in \mathbb{N}$ that

$$\mathbb{E}|F(\mathbf{k})|^{2p} \leq C^p p^p |\mathbf{k}|^{-p}.$$

Write $p_k = \lfloor \log |\mathbf{k}| \rfloor$. Then

$$\mathbb{E} \sum_{\mathbf{k} \in \varepsilon\mathbb{Z}^2 \setminus \{0\}} |\mathbf{k}|^{-3} \frac{|F(\mathbf{k})|^{2p_k}}{C^{p_k} p_k^{p_k} |\mathbf{k}|^{-p_k}} \leq \sum_{\mathbf{k} \in \varepsilon\mathbb{Z}^2 \setminus \{0\}} |\mathbf{k}|^{-3} < \infty.$$

This means that the summands tend to 0 almost surely as $|\mathbf{k}| \rightarrow \infty$ and so we can find a random constant $C_\omega > 0$ such that for all $\mathbf{k} \in \varepsilon\mathbb{Z}^2 \setminus \{0\}$ we have

$$|\mathbf{k}|^{-3} \frac{|F(\mathbf{k})|^{2p_k}}{C^{p_k} p_k^{p_k} |\mathbf{k}|^{-p_k}} \leq C_\omega.$$

Thus possibly making C_ω bigger, this yields

$$|F(\mathbf{k})| \leq C_\omega |\mathbf{k}|^{-1/2} \sqrt{\log |\mathbf{k}|}.$$

Now this holds for each $\mathbf{k} \in \varepsilon\mathbb{Z}^2 \setminus \{0\}$, so by the definition of $F(\mathbf{k})$ we have, whenever $\mathbf{k} = u(\cos \theta, \sin \theta) \in \varepsilon\mathbb{Z}^2 \setminus \{0\}$ with $\theta \in V_u$, that

$$|\hat{\mu}(\mathbf{k})| \leq C_\omega |\mathbf{k}|^{-1/2} \sqrt{\log |\mathbf{k}|}$$

as claimed. \square

3D. From lattices to \mathbb{R}^2 . We can now complete the proof of the main theorem. For this purpose, we need the following comparison lemma used by Kahane that allows us to pass from convergence on lattices for Fourier transforms to the whole space:

Lemma 3.11 [[Kahane 1985b](#), Lemma 1, page 252]. *Suppose τ is a measure on \mathbb{R}^2 with support in $(-1, 1)^2$. Suppose $\varphi, \psi : (0, \infty) \rightarrow (0, \infty)$ are decreasing as $t \rightarrow \infty$ with the doubling properties*

$$\varphi(t/2) = O(\varphi(t)) \quad \text{and} \quad \psi(t/2) = O(\psi(t)) \quad \text{as } t \rightarrow \infty.$$

If the Fourier transform of τ along the integer lattice \mathbb{Z}^2 satisfies

$$|\hat{\tau}(\mathbf{n})| = O\left(\frac{\varphi(|\mathbf{n}|)}{\psi(|\mathbf{n}|)}\right) \quad \text{as } |\mathbf{n}| \rightarrow \infty,$$

then

$$|\hat{\tau}(\xi)| = O\left(\frac{\varphi(|\xi|)}{\psi(|\xi|)}\right) \quad \text{as } |\xi| \rightarrow \infty.$$

Proof of Theorem 1.3. Combining Lemmas 3.7 and 3.3 we have that for any $\varepsilon > 0$, almost surely, there exists some random constant $C_\omega > 0$ such that for any $\mathbf{k} = u(\cos \theta, \sin \theta) \in \varepsilon \mathbb{Z}^2 \setminus \{0\}$ we have

$$|\hat{\mu}(\mathbf{k})| \leq C_\omega |\mathbf{k}|^{-1/2} \sqrt{\log |\mathbf{k}|}. \quad (3-4)$$

Define a measure τ_ε on \mathbb{R}^2 such that

$$\hat{\tau}_\varepsilon(\xi) := \hat{\mu}(\varepsilon \xi), \quad \xi \in \mathbb{R}^2.$$

By the almost sure continuity of W_t , we have that there exists a random constant $M_\omega > 0$ such that the diameter of the support of μ is at most M_ω almost surely. Taking an intersection of the events that (3-4) holds for $\varepsilon = 1/n$ over all $n \in \mathbb{N}$ allows us to find a random $\varepsilon = \varepsilon_\omega > 0$ such that μ is supported on a set of diameter strictly less than $1/\varepsilon$ and (3-4) holds almost surely with this ε . This guarantees that the measure τ_ε is supported on $(-1, 1)^2$ and so applying Lemma 3.11 with the measure $\tau = \tau_\varepsilon$ and the maps $\varphi(t) := \sqrt{\log t}$ and $\psi(t) := t^{1/2}$ gives the claim. \square

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Analytic torsion, dynamical zeta functions, and the Fried conjecture SHU SHEN	1
Existence theorems of the fractional Yamabe problem SEUNGHYEOK KIM, MONICA MUSSO and JUNCHENG WEI	75
On the Fourier analytic structure of the Brownian graph JONATHAN M. FRASER and TUOMAS SAHLSTEN	115
Nodal geometry, heat diffusion and Brownian motion BOGDAN GEORGIEV and MAYUKH MUKHERJEE	133
A normal form à la Moser for diffeomorphisms and a generalization of Rüssmann's translated curve theorem to higher dimensions JESSICA ELISA MASSETTI	149
Global results for eikonal Hamilton–Jacobi equations on networks ANTONIO SICONOLFI and ALFONSO SORRENTINO	171
High-frequency approximation of the interior Dirichlet-to-Neumann map and applications to the transmission eigenvalues GEORGI VODEV	213
Hardy–Littlewood inequalities on compact quantum groups of Kac type SANG-GYUN YOUN	237